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ON THE DETERMINATION OF OPTIMAL NODES AND  
WEIGHT FACTORS FOR THE SOLUTION OF COMPLETELY  
INCOHERENT SCATTERING PROBLEM

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**ABSTRACT.** For the numerical solution of the radiative transfer equation with complete redistribution of frequencies, it is necessary to select nodes and weight factors close to optimal, which provide the greatest proximity  $\delta$  between the exact and approximate solutions. The issue is reduced to a discrete boundary-value problem. We prove the theorem of existence and uniqueness of solution for above discrete boundary-value problem. We show that this solution imparts minimum to the function  $\delta$ . At the end of the work we give examples representing interest in radiative transfer theory of light and gamma quanta.

**Keywords:** nodes, weight factors, transfer equation, "shooting" method, existence of the solution, complete redistribution of frequencies.

## 1. THE STATEMENT OF PROBLEM

It is well known that the integral equation of radiative transfer of light and gamma quanta at complete redistribution of frequencies in one dimensional approximation has the form (see [1, 2]):

$$(1) \quad S(\tau) = g(\tau) + \int_0^{\infty} K(\tau - \tau') S(\tau') d\tau'.$$

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Here  $S(\tau)$  is the desired source function,  $g(\tau)$  is the free term,  $K(\tau)$  is the kernel of the integral equation

$$(2) \quad K(\tau) = \frac{\lambda}{2} A \int_{-\infty}^{+\infty} \alpha^2(x) e^{-\alpha(x)|\tau|} dx,$$

$\lambda$  is the quantum survival probability,  $A$  is the normalizing factor,  $\alpha(x)$  is an even function and represents absorption profile, and  $x$  is the dimensionless frequency.

For numerical solution of basic integral equation (1), we divide the interval  $(0, +\infty) : 0 = x_0 < x_1 < x_2 < x_3 < \dots < x_{n-1} < x_n < x_{n+1} = \infty$  and the kernel  $K(\tau)$  replace by a kernel  $\tilde{K}$  of the form

$$(3) \quad \begin{aligned} K(\tau) \approx \tilde{K}(\tau) &= \lambda A \sum_{m=0}^n e^{-\alpha_m |\tau|} \int_{x_m}^{x_{m+1}} \alpha^2(t) dt \\ &= \lambda A \sum_{m=0}^n a_m e^{-\alpha_m |\tau|} \end{aligned}$$

where

$$(4) \quad a_m = \int_{x_m}^{x_{m+1}} \alpha^2(t) dt; \quad \alpha_m = \alpha(x_m).$$

From (3), (4) it is easy to see that

$$(5) \quad \tilde{K}(\tau) \leq K(\tau).$$

We also consider the integral equation with kernel  $\tilde{K}(\tau)$

$$(6) \quad \tilde{S}(\tau) = g(\tau) + \int_0^\infty \tilde{K}(\tau - \tau') \tilde{S}(\tau') d\tau'.$$

We estimate the proximity between solutions  $S(\tau)$  and  $\tilde{S}(\tau)$  in space  $L_1^+(0, +\infty)$ . We have

$$(7) \quad \| S(\tau) - \tilde{S}(\tau) \|_{L_1^+} \leq \frac{\| g \|_{L_1^+}}{(1 - \lambda)(1 - \mu)} \delta,$$

where

$$(8) \quad \delta = \| K(\tau) - \tilde{K}(\tau) \|_{L_1^+},$$

$$(9) \quad \lambda = \int_0^\infty K(\tau) d\tau < 1, \quad \mu = \int_0^\infty \tilde{K}(\tau) d\tau < 1.$$

Thus the proximity of the kernels  $K$  and  $\tilde{K}$  provides proximity of the solutions to the basic equation (1) and corresponding "reduced" equation (6). In this connection, arises the problem to determine free parameters  $\alpha_m$  and  $a_m$  so that number  $\delta$  could reach its minimum at given fixed number of  $n$  nodes. It should be noted that when the kernel permits the representation

$$K(\tau) = \int_a^b e^{-s|\tau|} d\sigma(s), \quad 0 < a < b \leq +\infty$$

(coherent scattering) the above mentioned approach for the first time was suggested in work [3].

We have

$$(10) \quad \delta(x_1, x_2, x_3, \dots, x_n) = \int_0^\infty [K(\tau) - \tilde{K}(\tau)] d\tau = \lambda A \left[ 1 - \sum_{m=0}^n \frac{1}{\alpha_m} \int_{x_m}^{x_{m+1}} \alpha^2(t) dt \right].$$

Minimization of the function  $\delta$   $\left( \frac{\partial \delta}{\partial x_k} = 0 \right)$  yields the relation

$$(11) \quad -\frac{1}{\alpha_m^2} \alpha_m' \int_{x_m}^{x_{m+1}} \alpha^2(x) dx - \alpha_m + \frac{\alpha_m^2}{\alpha_{m-1}} = 0.$$

The boundary conditions are added to the relation (11)

$$(12) \quad x_0 = 0, \quad x_{n+1} = \infty.$$

The nodes  $x_k$  can be determined from equation

$$(13) \quad \int_{x_m}^{x_{m+1}} \alpha^2(x) dx = F(x_m, x_{m-1}),$$

where

$$(14) \quad F(x_m, x_{m-1}) = -\frac{\alpha_m^3}{\alpha_m'} \left( 1 - \frac{\alpha_m}{\alpha_{m-1}} \right).$$

We assume that  $x_0 = 0$ ,  $x_1 = t$  and from the equation (13) we determine  $x_1, x_2, x_3, \dots, x_k$  for a fixed number of  $n$ , at  $x_{n+1} = \infty$ . Here  $t$  is the parameter.

It should be noted that right hand side integral in equation (13) is monotonously increasing over  $x_k$ . If in finding the values of  $x_k$  ( $k < n$ ) we fail to determine the resulting  $x_{k+1}, x_{k+2}, \dots, x_n$ , i.e. equation (13) does not take place despite the increase of the integer upper limit, then the value of the parameter  $x_1 = t$  should be decreased. And the contrary, if the values of  $x_k$  ( $k > n$ ) are found, then the values of initial parameter  $x_1 = t$  should be increased. This process continues until all the nodes  $x_k$  ( $k = n$ ) are found.

Actually, from the boundary-value problem we proceed to a recurrent system. This approach is known as a "shooting" method, which often applied to the numerical solution of boundary-value problems (see [4]).

## 2. BASIC RESULT

Below we prove that nodes  $\{x_k\}_{k=0}^\infty$  being determined from (13) are unique for given fixed  $n$  and this solution imparts minimum to the function  $\delta$ .

**Theorem 1.** Let  $\alpha(0) = 1$ ,  $\alpha(\infty) = 0$ ,  $\alpha'(x) < 0$ ,  $x > 0$ . Let for arbitrary  $x_*$  there exists such number  $P_{x_*}$  that

$$(15) \quad |F(x, y)| \leq P_{x_*} |x - y|, \quad 0 \leq y \leq x \leq x_*,$$

where

$$(16) \quad F(x, y) = -\frac{\alpha^3(x)}{\alpha'(x)} \left[ 1 - \frac{\alpha(x)}{\alpha(y)} \right].$$

Then

a) for each  $n \in \mathbb{N}$  there exists unique solution  $(x_1, x_2, x_3, \dots, x_n)$  satisfying the equation

$$(17) \quad \int_{x_m}^{x_{m+1}} \alpha^2(x) dx = F(x_k, x_{k-1}), \quad x_0 = 0, \quad k = 0, 1, 2, \dots, n-1, \quad x_k = x_k(x_1),$$

b) discrete boundary-value problem (11), (12) has unique solution  $(x_1, x_2, x_3, \dots, x_n)$ , which imparts minimum to the function  $\delta$ .

*Proof.* We designate

$$(18) \quad H(x) = \int_x^\infty \alpha^2(t) dt - F(x, 0).$$

From conditions of theorem it follows that  $F(0, 0) = 0$ . Then

$$(19) \quad H(0) = \int_0^\infty \alpha^2(x) dx > 0.$$

Two cases are possible:

$$(20) \quad \text{i) } H(x) > 0 \text{ for } \forall x > 0 \quad \text{and} \quad \text{ii) } \exists \tilde{x} : H(\tilde{x}) = 0.$$

First we consider case i). We select  $x_1 > 0$ . Since  $\int_{x_1}^\infty \alpha^2(t) dt - F(x_1, 0) = H(x_1) > 0$ , then there exists unique number  $x_2$  such that

$$\int_{x_1}^{x_2} \alpha^2(t) dt = F(x_1, x_0).$$

Thus, we have found  $x_1, x_2, x_3, \dots, x_n$ . Then from obvious inequalities  $F(x_k, x_{k-1}) < F(x_k, 0)$  and  $H(x_k) > 0$  it follows that

$$\int_{x_k}^\infty \alpha^2(x) dx > F(x_k, x_{k-1}).$$

Hence there exists unique number  $x_{k+1}$ , such that

$$(21) \quad \int_{x_k}^{x_{k+1}} \alpha^2(x) dx = F(x_k, x_{k-1}).$$

The case ii). We denote

$$(22) \quad x_* = \min\{x : H(x) = 0\}, \quad m = \int_{x_*}^{\infty} \alpha^2(t) dt,$$

$$(23) \quad h = \max\{1, \alpha^{-2}(x_*)\}, \quad p = \max\{1, P_{x_*}\}.$$

We have

$$(24) \quad |x - y| \leq h \int_y^x \alpha^2(t) dt,$$

$$(25) \quad |F(x, y)| \leq p |x - y|, \quad 0 \leq y \leq x \leq x_*.$$

We select such a small point  $x_1$ , so that the both inequalities take place:

$$(26) \quad [1 + ph + p^2h^2 + \dots + p^{n-2}h^{n-2}]x_1 < x_*,$$

$$(27) \quad p^{n-2}h^{n-1}x_1 < m.$$

From (24) and (27) we have

$$(28) \quad F(x_1, x_0) \leq p |x_1 - x_0| = px_1 < m.$$

From (26) it follows that  $x_1 < x_*$ . Then  $\int_{x_1}^{\infty} \alpha^2(t) dt > m$  and therefore there exists  $x_2$ :

$$\int_{x_1}^{x_2} \alpha^2(t) dt = F(x_1, x_0),$$

besides,

$$(29) \quad |x_2 - x_1| \leq h \int_{x_1}^{x_2} \alpha^2(t) dt = hF(x_1, x_0) \leq phx_1.$$

From (29) it follows that

$$(30) \quad x_2 \leq (1 + p)hx_1, \quad |F(x_2, x_1)| \leq p |x_2 - x_1| \leq p^2hx_1 < m.$$

Thus, we have found  $x_1, x_2, x_3, \dots, x_n$ , ( $k \leq n - 2$ ) and

$$(31) \quad x_k \leq [1 + ph + \dots + (ph)^{k-1}]x_1 (< x_*),$$

$$(32) \quad |F(x_k, x_{k-1})| \leq p^k h^{k-1} x_1 (< m).$$

Then from (31) and (32) it follows that there exists unique  $x_{k+1}$  :

$$\int_{x_k}^{x_{k+1}} \alpha^2(t) dt = F(x_k, x_{k-1}).$$

Taking into consideration (31), (32), (24), we obtain

$$x_{k+1} \leq [1 + ph + \dots + (ph)^k]x_1 (< x_*),$$

$$|F(x_{k+1}, x_k)| \leq p^{k+1} h^k x_1 (< m).$$

Thus by induction we show that if  $k \leq n - 2$ , then there exists unique  $x_{k+1}$  :

$$(33) \quad \int_{x_k}^{x_{k+1}} \alpha^2(t)dt = F(x_k, x_{k-1})$$

and the conditions (31), (32) are fulfilled. Imposing  $k = n - 1$  and using (24), (25), (31), (32) we conclude that

$$\exists x_n : \int_{x_{n-1}}^{x_n} \alpha^2(t)dt = F(x_{n-1}, x_{n-2}).$$

The theorem is proved. □

### 3. EXAMPLES

It is easy to check that functions  $\alpha(x) = e^{-x^2}$  (Doppler profile) and  $\alpha(x) = \frac{1}{1+x^2}$  (Lorentz profile), representing practical interest in radiative transfer theory of light and gamma quanta (see [1]), satisfy conditions of theorem .

Lets verify the condition (15).

In the case of the Doppler profile  $\alpha(x) = e^{-x^2}$ ,  $A = \frac{1}{\sqrt{\pi}}$  we have

$$| F(x, y) | = \frac{e^{-2x^2}}{2x} (1 - e^{-(x^2 - y^2)}) \leq \frac{e^{-2x^2}}{2x} (x^2 - y^2) \leq |x - y|, \quad p = 1, \quad 0 \leq y \leq x \leq x_*$$

The relation (13) takes the form

$$\int_{x_m}^{x_{m+1}} \alpha^2(t)dt = \frac{e^{-2x_m^2}}{2x_m} \left[ 1 - \frac{e^{-x_m^2}}{e^{-x_{m-1}^2}} \right], \quad a_m = \frac{1}{\sqrt{\pi}} \int_{x_m}^{x_{m+1}} e^{-t^2} dt.$$

In case of Lorentz profile  $\alpha(x) = \frac{1}{1+x^2}$ ,  $A = \frac{1}{\pi}$ ,

$$| F(x, y) | = \frac{1}{2x(1+x^2)} \left( \frac{x^2 - y^2}{1+x^2} \right) \leq \frac{x-y}{1+x^2} \leq \frac{1}{2}(x - y), \quad p = 1, \quad 0 \leq y \leq x \leq x_*$$

The relation (13) takes the form.

$$\frac{x_{m+1}}{x_{m+1}^2 + 1} + \arctg x_{m+1} = \frac{x_m}{x_m^2 + 1} + \arctg x_m + \frac{x_m^2 - x_{m-1}^2}{x_m(x_m^2 + 1)^2},$$

$$a_m = \frac{1}{\pi} \left[ \frac{x_{m+1}}{x_{m+1}^2 + 1} - \frac{x_m}{x_m^2 + 1} + \arctg x_{m+1} - \arctg x_m \right].$$

### REFERENCES

[1] V.V. Ivanov, *Perenos Izluchenii i Spectra Nebesnikh Tel.*, M. Nauka, 1969, 472p.  
 [2] A.Kh. Khachatryan, *On One Numerical Method of Solution of Coherent and Incoherent Scattering Problem*, *Astrophysics and Space Science*, **215** (1994), 17-35.  
 [3] N.B.Engibaryan, E.A.Melkonyan, *On Discrete Ordinate Method*, *Doclady AN SSSR*, **292** (1987), 322-326. (In russian) *Zbl* 0641.65095  
 [4] J.Stoer, R.Bulirsch, *Introduction to Numerical Analysis*, New York, Springer-Verlag, 1980. *Zbl* 0423.65002

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